Phase-modulator-based optoelectronic oscillator for generating short optical pulse and microwave signal

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Abstract. A phase-modulator-based optoelectronic oscillator (OEO) is utilized to implement a self-starting optical pulse and microwave signal source. This system is able to simultaneously generate a 9.8-GHz optical pulse stream with 141-fs (over 100 Hz to 1 MHz) timing jitter and 6.5-ps pulse width, along with a high spectral purity electrical signal that is locked to the repetition rate of the optical pulses. The measured phase noise is −112 dBc/Hz at 10 kHz away from carrier, and the side modes are perfectly suppressed. ⊚ 2007 Society of Photo-Optical Instrumentation Engineers. [DOI: 10.1117/1.2775962]

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1 Introduction

The optoelectronic oscillator (OEO) has attracted many researchers because it is a very useful optical and electrical microwave signal source. Recently, a few significant works for extending the concept of OEO were reported, for instance, an OEO performs as an optical pulse source without an external microwave-driving source. But the reported pulse width is around 15 to 20 ps (at 10 GHz), thin in their applications in high-speed optical time division multiplexing (OTDM) systems.

Compared with these schemes, an electro-optic phase modulator (PM)-based OEO is another effective way to perform the same function. ^{6,7} When continuous-wave (cw) light is injected in an electro-optic PM, it is deeply chirped. After an optical synthesizer or group-delay-dispersion (GDD) circuit, the chirped cw light is compressed and trains of short pulses are generated. Enhancement of linear chirp is the key to generating short pulses. Furthermore, performing amplitude modulation (AM) before PM is effective to eliminate the residual optical power in the midpoint between adjacent pulses after compression. ⁶

We present a scheme for a phase-modulator-based dualloop OEO, which can simultaneously produce ultra-lowjitter short optical pulse streams with high repetition rate capabilities, and corresponding high spectral purity and low phase noise electrical microwave signals. A joint of a large-

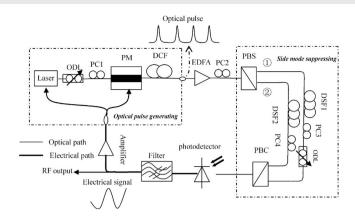


Fig. 1 The experimental setup.

signal directly modulated laser (DML) and a phase modulator is for chirp enhancing and pulse shaping. An optical domain combined dual-loop configuration is employed for side-mode suppression without adding electrical active devices. The experimental setup is depicted in Fig. 1.

2 Principle of Operation

As shown in Fig. 1, the system includes a directly modulated distributed feedback (DFB) laser, a LiNbO3 phase modulator with V_{π} of 3.5 V, an erbium-doped fiber amplifier (EDFA), a high-speed PIN photodetector, and an rf amplifier centered at 9.8 GHz with 150-MHz bandwidth. The electrical filter is not used due to the signal frequency already limited by the rf amplifier. The directly modulated light is injected into the PM after adjustment of the optical delay line (ODL) to achieve the maximum chirping rate. Polarization controller 1 (PC1) is used to obtain the best modulation. The EDFA is used for compensating the insertion loss in the system. PC2 controls the power distribution in optical path1 and path2, which are 10 and 2 km, respectively. PC3 and PC4 tune the power distribution and correct the polarization variation caused by dispersion shift fiber (DSF). Once the oscillation threshold is reached, a stable output is achieved.

The output of large-signal DML can be approximately expressed as:

$$E(t) = E_0 \exp\left(-\frac{t^2}{2t_0^2}\right) \exp\left[i(\omega_0 t - \varphi(t))\right],\tag{1}$$

where t_0 is the pulse width, which is defined as half the width of the optical pulse when amplitude decreases to 1/e of peak. ω_0 is the angular optical frequency and $\varphi(t)$ is the phase variation depending on time t. When this light field is phase modulated by a sinusoidal signal of frequency f_m with a time delay (Δt) between DML and PM, we obtain 7

$$E_{PM}(t) = E_0 \exp\left(-\frac{t^2}{2t_0^2}\right) \exp\{i[\omega_0 t - \varphi(t) - \Delta\theta \sin 2\pi f_m(t + \Delta t)]\}, \tag{2}$$

where $\Delta \theta$ is the modulation index. Suppose $f_m \cdot \Delta t = k$ (k: integer), the instantaneous frequency of the modulated light is then given by

$$\nu(t) = \omega_0/2\pi - \partial\varphi(t)/2\pi\,\partial\,t - \Delta\,\theta f_m\cos(2\pi f_m t). \tag{3}$$

The slope of the second term is nearly linear when it closes to the center frequency of light. In addition, near $f_m t = n/2 + 1/4$ (n: integer), the third term gives almost a linear optical frequency chirp. Then, the extremum chirping rates can be written as

$$(\partial \nu / \partial t)_{\text{max}} = -\partial^2 \varphi(t) / 2\pi \partial t^2 \pm 2\pi f_m^2 \Delta \theta. \tag{4}$$

Because dispersion compensation fiber (DCF) can only compensate the linear negative chirps, the timing between DML and PM is controlled by an adjustable optical delay line, such that the negative chirp of the DML was gated in the PM, generating optical pulses with a flat optical frequency spectrum. The maximum chirping rate is rewritten as follows:

$$(\partial \nu / \partial t)_{-\text{max}} = -\frac{1}{2\pi} |\partial^2 \varphi(t) / \partial t^2| - 2\pi f_m^2 \Delta \theta.$$
 (5)

Obviously, the frequency chirp is larger than that performed by DML or PM alone. After the compression with the optimum GDD, a shorter optical pulse stream is generated. For PM, a larger modulation depth is effective to gain shorter pulse, which is only limited by the endurance of the E/O crystal. Notice that the DML is not only for enhancing chirp but also for eliminating residual optical power (corresponding to the function of AM before PM), although still a minor part of energy is dispersed as the pedestal, and nonlinear chirps are formed as small peaks among pluses.

To suppress side modes, a polarization beamsplitter (PBS) and polarization beam combiner (PBC) are employed to compose an optical domain combined dual-loop structure. Little close-to-carrier noise is introduced, and the synthetical noise performance is greatly improved.

Notice that DSF is used in this experiment to prevent extra GDD. It is well known that the dispersion of DCF is equal in magnitude and opposite in sign to that of the single-mode fiber (SMF). As a result, the pulses compressed by DCF are rebroadened if SMF is used. In addition, the transmission of short pulses in DSF will excite self-phase modulation, which provides the potential for further compression.

3 Experimental Results and Conclusion

In this experiment, breaking either optical path between PBS and PBC, it is a completed single loop OEO that can freely run when enough optical power and amplifier gain is supplied. For an electrical signal, with path 1 broken, the measured side-mode suppression ratio is 55 dB and mode spacing is 93.3 kHz [see Fig. 2(a)]. With path 2 broken, the measured side-mode suppression ratio is 45 dB and mode spacing is 19.8 kHz [see Fig. 2(b)]. In the dual-loop case, the side modes are perfectly suppressed [see Fig. 2(c)] and the measured phase noise is -112 dBc/Hz at 10 kHz [see Fig. 2(d)]. To evaluate the timing jitter (σ) of the optical pulse stream, we integrate the spectral noise-power density from 100-Hz to 10-MHz offset using the following equation:

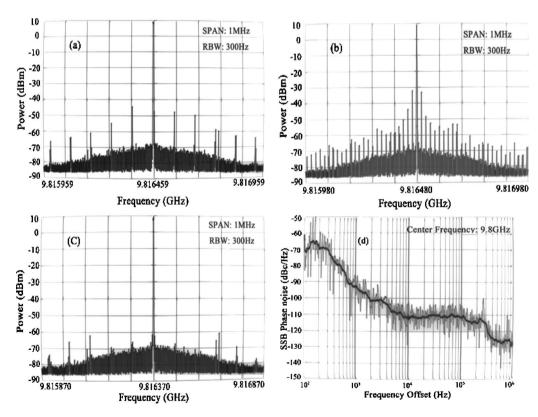
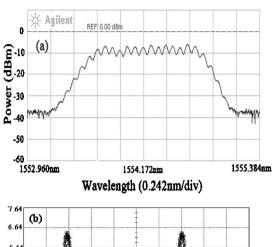
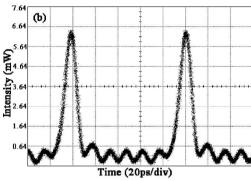


Fig. 2 Experiment results of the electrical signal: (a) the RF spectrum of signal loop OEO with 2-km loop length at SPAN 1 MHz, RBW 300 Hz; (b) rf spectrum of single loop OEO with 10-km loop length at SPAN 1 MHz, RBW 300 Hz; (c) rf spectrum of dual-loop OEO at SPAN 1 MHz, RBW 300 Hz; and (d) SSB phase noise spectrum of the PM-based dual-loop OEO.





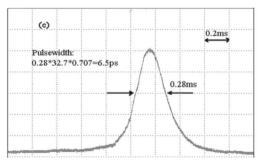


Fig. 3 Optical pulse characterization: (a) optical spectrum with 0.06-nm resolution, (b) time domain measurement of the 9.8-G optical pulse, and (c) autocorrelation trace.

$$\sigma = \left[2 \int_{f_{\min}}^{f_{\max}} L(f) df \right]^{1/2} / 2\pi f_{\text{osc}},$$

where the $f_{\rm osc}$ is the oscillation frequency. The yielding result of timing jitter is 141 fs. The phase noise and timing jitter performance is mainly influenced by MDL and EDFA, which introduce amplitude fluctuation and noise, as well as timing jitter. Furthermore, the random fiber-length fluctuations caused by environment effects are other significant factors.

Figure 3(a) shows the optical spectrum of the pulse stream. The full-width at half-maximum (FWHM) is measured to be 0.830 nm, which is about 0.3 nm wider than that generated by large-signal DML alone. In this case, the

modulation voltage is 3.54 V (corresponding to the modulation depth of 1π) on PM and 22.3 dBm on DML with a bias of 50 mA. The sampling oscilloscope trace for the optical pulse stream is given in Fig. 3(b), where the pulse width is 6.5 ps measured by an autocorrelator [see Fig. 3(c)]. It is also clearly shown in Fig. 3(b) that the optical energy at the tail far from the pulse peak is almost perfectly eliminated. The bandwidth product of $\Delta\nu\Delta t\sim0.65$ is not transform limited for an assumed Gaussian pulse shape, because the dispersion magnitude (-21.68 ps/nm) of DCF is not optimum and the modulated light includes uncompressible nonlinear chirps. Nevertheless, this result is better than that generated by DML or PM alone.

In conclusion, we have demonstrated a self-starting phase-modulator-based OEO for generating low-timing jitter short optical pulses and low phase noise rf signals. Pulse shaping and side-mode suppression techniques are employed in the cavity of OEO. The output is 6.5-ps-wide pulses at 9.8-GHz repetition rate with 141-fs timing jitter. A low phase noise, high-quality, 9.8-GHz electrical signal is simultaneously generated that can serve as a synchronized reference in applications of this pulse source. The phase noise and timing jitter performance can be improved by controlling environment effects and using a high performance light source. In addition, the pulse rate can be directly scaled to 40 GHz or higher by changing some opto-electronic components in the system.

Acknowledgments

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